PAUL SCHERRER INSTITUT



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## **Implicit Solvent Methods and DFT+U**

CP2K User Tutorial: "Computational Spectroscopy", Uni Paderborn, 28 August 2018



## DFT+U



## Local density approximation (LDA)

- Plain DFT (LDA, GGA) performs unexpectedly well for many solid state systems
- But:
  - LDA provides a poor description of the Coulomb interaction
  - The LDA functional is derived from a weakly correlated reference system, the jellium model (uniform or homogeneous electron gas)
  - Generalised Gradient Approximation (GGA) does not overcome these shortcomings as it is still a "local" approximation by adding only gradient information
  - Fortuitous error cancellation obscures the deficiencies
  - LDA fails badly for strongly correlated systems like transition metal (TM) oxides and heavy fermion materials with partially filled 4f and 5f orbitals



- Hubbard bands form as an effect of the electronic correlation:
  - LDA bands split into two sets of bands separated by a local Coulomb repulsion U and band width  $W \rightarrow$  Mott insulator





- Though it is possible to calculate the Hubbard parameter *U* using linear response theory\*, it is usually treated as an empirical fitting parameter
- Fixed *U* parameter does not adapt to the chemical environment (e.g. change of oxidation state)
- Occurrence of meta-stable states for heavy fermion materials like CeO<sub>2</sub>, UO<sub>2</sub> etc.
   ⇒ major obstacle technically
- On-going development, e.g. dynamical mean field theory (LDA+DMFT)
  - However, such methods are computationally very expensive
  - Accessible model system sizes are quite limited currently (≈ 30 atoms per cell)
  - Analytical energy gradients, i.e. atomic forces, are not straightforwardly available

\*Cococcioni and Gironcoli, Phys. Rev. B **71**, 035105



- Hubbard correction term is added to the DFT energy functional as an atomic onsite term (V. I. Anisimov et al., Phys. Rev. B 44, 943 (1991))
- Explicit correction that acts as an energy penalty function, e.g. Dudarev\*:

$$E_{\rm U} = \frac{\sigma_{\rm eff}}{2} \sum_{\sigma,I} {\rm Tr}[\boldsymbol{n}^{\sigma,I} (\boldsymbol{1} - \boldsymbol{n}^{\sigma,I})]$$
 with atom *I* and spin  $\sigma$ 





- Orbital based methods:
  - Chirgwin and Coulson:  $P = \frac{1}{2}(DS + SD)$
  - Mulliken: P = D S
  - Löwdin:  $P = S^{1/2} D S^{1/2}$

#### with density matrix **D**, overlap matrix **S**, and occupation matrix **P**

- Partitioning of the electronic density:
  - Bader charges: Gradient of the electronic charge density (zero-flux surface)
  - Hirshfeld charges: Difference to unrelaxed atomic densities (pro-density)
- Potential derived charges methods
- Moment derived charges



## Tackling the meta-stable state problem

- Occupation matrix control (OMC)
   B. Dorado et al., Phys. Rev. B 79, 235125 (2009)
- U ramping method B. Meredig et al., Phys. Rev. B 82, 195128 (2010)
- Quasi-annealing (QA) method H. Y. Geng et al., Phys. Rev. B 82, 094106 (2010)
- Controlled symmetry reduction (CSR) method D. Gryaznov et al., Phys. Chem. Chem. Phys. **14**, 4482 (2012)
- f occupation smearing and U ramping (FOUR) J. Rabone and M. Krack, Comput. Mat. Sci. **71**, 157 (2013)
- Local electronic minima inhibition by averaging occupations (LEMIAO)
   J. Rabone and M. Krack, Comput. Mat. Sci. 71, 157 (2013)







Force evaluation section

• Antiferromagnetic ground state below 30.4 K:

```
&FORCE_EVAL
METHOD Quickstep
 STRESS_TENSOR analytical
 &DFT
  CHARGE 0
  LSD
  MULTIPLICITY 1
  PLUS U METHOD Mulliken
  • • •
 &END DFT
    •
```

&END FORCE\_EVAL



- All-electron configuration of uranium: 92 electrons  $1s^2 2s^2 2p^6 3s^2 3p^6 3d^{10} 4s^2 4p^6 4d^{10} 4f^{14} 5s^2 5p^6 5d^{10} 6s^2 6p^6 5f^3 6d^1 7s^2$
- 86 core electrons:  $1s^2 2s^2 2p^6 3s^2 3p^6 3d^{10} 4s^2 4p^6 4d^{10} 4f^{14} 5s^2 5p^6 5d^{10} 6s^2 6p^6$
- 6 valence electrons:

 $5f^3 6d^1 7s^2$ 



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- 6 valence electrons:

 $5f^3 6d^1 7s^2$ 



- All-electron configuration of oxygen: 8 electrons  $1s^2 2s^2 2p^4$
- 2 core electrons:

 $1s^{2}$ 

• 6 valence electrons:

 $2s^2 2p^4$ 



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 $1s^{2}$ 

• 6 valence electrons:

 $2s^2 2p^4$ 



- All-electron configuration of oxygen: 8 electrons  $1s^2 2s^2 2p^4$
- Pseudo atom configuration of oxygen: 2 + 6 electrons [He]  $2s^2 2p^4$



- All-electron configuration of uranium: 92 electrons  $1s^2 2s^2 2p^6 3s^2 3p^6 3d^{10} 4s^2 4p^6 4d^{10} 4f^{14} 5s^2 5p^6 5d^{10} 6s^2 6p^6$  $7s^2 5f^3 6d^1$
- Large-core pseudo atom configuration of uranium: 86 + 6 electrons [Rn]  $7s^2 5f^3 6d^1$
- Medium-core pseudo atom configuration of uranium: 78 + 14 electrons [Xe  $4f^{14} 5d^{10}$ ]  $6s^2 6p^6 7s^2 5f^3 6d^1$
- Small-core pseudo atom configuration of uranium: 60 + 32 electrons [Kr  $4d^{10} 4f^{14}$ ]  $5s^2 5p^6 5d^{10} 6s^2 6p^6 7s^2 5f^3 6d^1$



**DFT+U** section in the atomic kind section

```
&KIND O
BASIS_SET DZVP-MOLOPT-SR-GTH-q6
POTENTIAL GTH-PBE-q6
&BS
  • • •
&END BS
```

```
! Not needed for O
&END KIND
```



**DFT+U** section in the atomic kind section

```
&KIND Ua
BASIS_SET DZVP-MOLOPT-GTH-q14
 ELEMENT U
POTENTIAL GTH-PBE-q14
&BS
  • • •
&END BS
&DFT PLUS U on
  L 3
 U_MINUS_J [eV] 2.00
  &ENFORCE OCCUPATION on/off
   . . .
  &END ENFORCE_OCCUPATION
&END DFT PLUS U
&END KIND
```



• Electronic configuration of U<sup>4+</sup>:

 $\begin{array}{l} 1s^2\ 2s^2\ 2p^6\ 3s^2\ 3p^6\ 3d^{10}\ 4s^2\ 4p^6\ 4d^{10}\ 4f^{14}\ 5s^2\ 5p^6\ 5d^{10}\\ 6s^2\ 6p^6\ 5f^2\end{array}$ 

• Medium-core pseudo atom configuration of U<sup>4+</sup>:  $[Xe 4f^{14} 5d^{10}] 6s^2 6p^6 5f^2$ 

• Electronic configuration of O<sup>2-</sup>:

$$1s^2 2s^2 2p^6$$

• Pseudo atom configuration of O<sup>2-</sup>:

[He] 
$$2s^2 2p^6$$



• Set up the (on-site) atomic orbital occupations for O<sup>2-</sup>:

$$0 \xrightarrow{+2e}{\longrightarrow} 0^{2-}$$
$$2p^4 \xrightarrow{+2e}{\longrightarrow} 2p^6$$

• BS section input is not processed quite intuitively:

```
&BS
    &ALPHA
    N   2 ! 2
    L   1 ! p
    NEL +2 ! (4 + 2)/2 = 3 alpha 2p electrons
    &END ALPHA
    &BETA
    N   2 ! 2
    L   1 ! p
    NEL +2 ! (4 + 2)/2 = 3 beta 2p electrons
    &END BETA
&END BETA
&END BS
```



• Set up the (on-site) atomic orbital occupations for U<sup>4+</sup>:

$$U \xrightarrow{-4e} U^{4+}$$
$$5f_{\alpha}^{3} 6d_{\alpha}^{1} 7s^{2} \xrightarrow{-4e} 5f_{\alpha}^{2}$$

 On-site triplet state for U<sup>4+</sup> for the spin-up (alpha) uranium atoms Ua and swap &ALPHA and &BETA sections for the spin-down (beta) Ub kind (not shown):

```
&BS
&AI PHA
         5
 Ν
              6
        3 2
  Т
                  0
       +1
             -1 -2 ! (3 + 1)/2 = 2 alpha 5f electrons
 NEL
&END ALPHA
&BETA
 Ν
        5
              6
                  7
             2
        3
                  0
  L
                 -2 ! remove all beta valence electrons
             -1
 NFI
        - 3
&END BETA
&END BS
```



• Two alpha electrons in seven 5f orbitals:  $\binom{7}{2} = 21$  combinations

	<i>f</i> <sub>-3</sub>	<i>f</i> <sub>-2</sub>	$f_{-1}$	$f_0$	$f_{+1}$	<i>f</i> <sub>+2</sub>	<i>f</i> <sub>+3</sub>
1	↑	1					
2	ſ		1				
3	↑			↑			
4	↑				↑		
5	Ť					↑	
6							ſ
7		1	1				
8		ſ		↑			
9		1			1		
10		1				1	

_	_				_		
	$f_{-3}$	$f_{-2}$	$f_{-1}$	$f_0$	$f_{+1}$	<i>f</i> <sub>+2</sub>	<i>f</i> +3
11		Ť					↑
12			1	1			
13			ſ		↑		
14			ſ			↑	
15			1				↑
16				↑	↑		
17				↑		↑	
18				↑			ſ
19					1	1	
20					1		↑
21						↑	↑



## Print DFT+U specific information

• Print occupation of each atom with a *U* value greater zero:

```
&GLOBAL
 PRINT_LEVEL medium
 . . .
&END GLOBAL
&FORCE_EVAL
 &DFT
  • • •
  &PRINT
   &PLUS_U on
    &EACH
     QS SCF 1
    &END EACH
   &END PLUS U
  &END PRINT
  . . .
 &END DFT
  . .
&END FORCE EVAL
```



### Print DFT+U specific information

• Use a tiny U value to trigger printout for the U = 0 case:

```
&DFT_PLUS_U
...
U_MINUS_J [eV] 2.00E-20
...
&END DFT_PLUS_U
```



#### **Occupation control**

• A specific (initial) orbital occupation can be enforced:

```
&DFT_PLUS_U
...
&ENFORCE_OCCUPATION on/off
EPS_SCF 1.0E-5 ! Enforce until a certain SCF convergence is reached
MAX_SCF 20 ! Enforce occupation for first 20 SCF iterations
ORBITALS -3 -2 -1 +0 +1 +2 +3 ! Smear f electrons over all f orbitals
SMEAR on
&END ENFORCE_OCCUPATION
...
&END DFT_PLUS_U
```

• Alternatively, define and enforce a specific orbital occupation pattern, e.g.

```
ORBITALS -3 -1
```



#### **Example: XAFS spectra simulation**





# **Implicit Solvent Methods**



## Modelling of solvation effects

- Vacuum
- Gas phase
- Explicit solvent with solute





## Modelling of solvation effects

- Vacuum
- Gas phase
- Explicit solvent with solute
- Implicit solvent with solute





**Continuum solvation models** 

- Polarisable continuum model (PCM, Tomasi et al.)
- Conductor-like screening model (COSMO, Klamt et al.)
- Smooth dielectric models are needed for molecular dynamics (MD) simulations



## Modelling of solvation effects

- Vacuum
- Gas phase
- Explicit solvent with solute
- Implicit solvent with solute





## Modelling of solvation effects

- Vacuum
- Gas phase
- Explicit solvent with solute
- Implicit solvent with solute
- Smoothed solute-solvent transition





- Explicit solute-solvent interactions are missing
  - Pros:
    - Shortcomings of standard DFT in the description of van der Waals and hydrogen bonds w.r.t. the solvent do not matter
    - Shorter sampling times needed to obtain meaningful thermodynamic averages
  - Cons:
    - No detailed description of hydrogen bonds (network)
    - Potential (chemical) reactions between solvent and solute are excluded a priori
- No explicit solvent atoms
  - Pros:
    - Significant reduction in problem size especially for localised basis set methods
    - More efficient sampling
  - Cons:
    - Additional (nested) convergence cycle for polarisation charge required
    - No information about solvation shells (detailed feedback of the solvent)
    - Potentially more noisy forces due to finite differences approach



**Smoothed dielectric function** 

• Dielectric as a smoothed self-consistent function of the electronic density:

$$\epsilon(\mathbf{r}) \equiv \epsilon \left[ \rho^{\text{elec}}(\mathbf{r}) \right] = \begin{cases} 1 & \text{large } \rho^{\text{elec}} \\ \epsilon_0 & \rho^{\text{elec}} \to 0 \end{cases}$$

J.-L. Fattebert and F. Gygi, J. Comput. Chem. 23, 662 (2002)
J.-L. Fattebert and F. Gygi, Int. J. Quantum Chem. 93, 139 (2003)
O. Andreussi et al., J. Chem. Phys. 136, 064102 (2012)



• Adding a dielectric medium to the system

$$\epsilon(\mathbf{r}) \equiv \epsilon[\rho^{\text{elec}}(\mathbf{r})]$$

 $\nabla^{2}\phi^{\text{tot}}(\boldsymbol{r}) = -4\pi\rho^{\text{solute}}(\boldsymbol{r}) \quad \text{with } \rho^{\text{solute}}(\boldsymbol{r}) = \rho^{\text{ions}}(\boldsymbol{r}) + \rho^{\text{elec}}(\boldsymbol{r})$  $\nabla \cdot \boldsymbol{\epsilon} [\rho^{\text{elec}}(\boldsymbol{r})]\phi^{\text{tot}}(\boldsymbol{r}) = -4\pi\rho^{\text{solute}}(\boldsymbol{r})$  $\nabla \cdot \boldsymbol{E}(\mathbf{r}) = 4\pi\rho^{\text{solute}}(\boldsymbol{r}) - 4\pi \cdot \boldsymbol{P}(\boldsymbol{r})$  $\rho^{\text{pol}}(\boldsymbol{r}) = -\nabla \cdot \boldsymbol{P}(\boldsymbol{r}) = \nabla \cdot \left(\frac{\epsilon(\rho^{\text{elec}}(\boldsymbol{r})) - 1}{4\pi}\nabla\phi^{\text{tot}}(\boldsymbol{r})\right)$ 

• Finally a vacuum-like Poisson problem is recovered

$$\nabla^2 \phi^{\text{tot}}(\boldsymbol{r}) = -4\pi \big( \rho^{\text{solute}}(\boldsymbol{r}) + \rho^{\text{pol}}(\boldsymbol{r}) \big)$$



Self-consistent Continuum Solvation (SCCS)

• Vacuum-like Poisson problem is recovered

$$\nabla^2 \phi^{\text{tot}}(\boldsymbol{r}) = -4\pi \big( \rho^{\text{solute}}(\boldsymbol{r}) + \rho^{\text{pol}}(\boldsymbol{r}) \big)$$

• Energy term

$$E^{\rm el} = E^{\rm solute} + E^{\rm pol}$$



• Solvation free energy

$$\Delta G^{\rm sol} = \Delta G^{\rm el} + G^{\rm rep} + G^{\rm dis} + G^{\rm cav} + \Delta G^{\rm tm} + P\Delta V$$

• Electrostatic contribution:

$$\Delta G^{\rm el} = G^{\rm el} - G^{\rm 0}$$

with the energy  $G^0$  of the solute in vacuum

• Repulsion term\*:

 $G^{\mathrm{rep}} = \boldsymbol{\alpha} S$ 

where S is the (quantum) surface of the cavity

• Dispersion term\*:

$$G^{\mathrm{dis}} = \beta V$$

where V is the (quantum) volume of the cavity

```
*Scherlis et al., J. Chem. Phys. 124, 074103 (2006)
```



• Cavitation term\*:

$$G^{\mathrm{cav}} = \gamma S$$

where S is the (quantum) surface of the cavity

- Thermal motion term  $G^{tm}$  and the volume change term  $P\Delta V$  are often ignored
- Collecting all terms results in an approximation for the solation free energy  $\Delta G^{\text{sol}} \approx \Delta G^{\text{el}}(\epsilon_0, \rho_{\min}, \rho_{\max}) + (\alpha + \gamma) S + \beta V$

- Quantum volume V and surface S:
  - M. Cococcioni et al., Phys. Rev. Lett. 94, 145501 (2005)

\*Scherlis et al., J. Chem. Phys. 124, 074103 (2006)



**SCCS** input section

```
&SCCS on/off
ALPHA [N*m^-1] 0.0
 BETA [kbar] 0.0
 DELTA RHO 2.0E-5
 DERIVATIVE_METHOD cd3/cd5/cd7/fft
 DIELECTRIC CONSTANT 78.36
 EPS_SCCS 1.0E-6
 GAMMA [mN/m] 0.0
@IF ${OT}
 EPS SCF 0.03
@ENDIF
@IF ${TD}
 EPS SCF 0.3
@ENDIF
 MAX ITER 100
 METHOD Andreussi/Fattebert-Gygi
 MIXING 0.6
 • • •
&END SCCS
```



• Fattebert-Gygi

$$\epsilon[\rho^{\text{elec}}(\boldsymbol{r})] = 1 + \frac{\epsilon_0 - 1}{2} \left( 1 + \frac{1 - \left(\rho^{\text{elec}}/\rho_0\right)^{2\beta}}{1 + \left(\rho^{\text{elec}}/\rho_0\right)^{2\beta}} \right)$$

• Andreussi et al.

$$\epsilon[\rho^{\text{elec}}(\boldsymbol{r})] = \begin{cases} 1 & \rho^{\text{elec}} > \rho_{\text{max}} \\ \exp\left(t(\ln\rho^{\text{elec}})\right) & \rho_{\text{min}} < \rho^{\text{elec}} < \rho_{\text{max}} \\ \epsilon_0 & \rho^{\text{elec}} < \rho_{\text{min}} \end{cases}$$

$$t(x) = \frac{\ln \epsilon_0}{2\pi} \left[ 2\pi \frac{\ln \rho_{\max} - x}{\ln \rho_{\max} - \ln \rho_{\min}} - \sin \left( 2\pi \frac{\ln \rho_{\max} - x}{\ln \rho_{\max} - \ln \rho_{\min}} \right) \right]$$



## Smoothing methods for the dielectric function

```
&SCCS on/off
 • • •
 DIELECTRIC_CONSTANT 78.36
 METHOD Andreussi/Fattebert-Gygi
 • • •
 &ANDREUSSI
  RHO_MAX 0.001
  RHO MIN 0.0001
 &END ANDREUSSI
 . . .
 &FATTEBERT-GYGI
  BETA 1.3
  RHO ZERO 0.0004
 &END FATTEBERT-GYGI
 . . .
```

&END SCCS



**Questions or comments?** 

